Comment on "Josephson Current through a Nanoscale Magnetic Quantum Dot"

In a recent work [1], Siano and Egger (SE) studied the Josephson current through a quantum dot in the Kondo regime using the quantum Monte Carlo (QMC) method. Several of their results were unusual, and inconsistent with those from the numerical renormalization group (NRG) calculations[2, 3] among other previous studies. Those results in Ref.[1] are not reliable for the following two reasons: (i) The definition of the Kondo temperature was wrong; (ii) There were substantial finite-temperature effects.

We first clarify point (i). The normal-state Kondo temperature [4, 5] in the absence of superconductivity provides one of the most significant energy scales of the system. SE defined the Kondo temperature as

$$T_K^{\rm SE} = \exp[\pi \epsilon_0 (\epsilon_0 + U) / \Gamma_{\rm SE} U] \sqrt{\Gamma_{\rm SE} U} / 2$$
 (1)

with $\Gamma_{\rm SE} = 2\pi \rho_0 |t|^2$, where $|t|^2$ denotes the coupling to one lead and ρ_0 the density of states (DOS) at the Fermi level. In Ref.[2] we defined it as

$$T_K = \exp[\pi \epsilon_0 (\epsilon_0 + U)/2\Gamma U] \sqrt{\Gamma U/2}$$
 (2)

with $\Gamma = 2\pi N_0 |V|^2$, where $|V|^2$ denotes the coupling to one lead and N_0 the DOS at the Fermi level per spin [the factor 2 in the coupling comes from the two leads]. It is important to clarify the difference between the two definitions since different definitions of T_K result in significantly different scaling behaviors of physical quantities. We note that both forms, Eqs. (1) and (2), appear in the literature. However, in Eq. (1) $\Gamma_{\rm SE}$ should be the full width at half maximum of the single particle level of the noninteracting dot [6], whereas in Eq. (2) Γ should be the half width at half maximum (HWHM) of the single particle level. To see the precise meaning of Γ_{SE} , let us take the limit $\Delta = 0$ and U = 0 in the local Green's function (GF) in Eq. (6) in Ref. [1]. Going over to the retarded GF, we find $G_R^{-1} \sim E + i \Gamma_{SE}$, which yields the spectral function

$$A(E) = -\frac{1}{\pi} \text{Im} G_R = \frac{1}{\pi} \frac{\Gamma_{SE}}{E^2 + \Gamma_{SE}^2}.$$
 (3)

Therefore, $\Gamma_{\rm SE}$ corresponds to HWHM. Namely, the two hybridization $\Gamma_{\rm SE}$ and Γ are the same. Therefore, the two Kondo temperatures in Eqs. (1) and (2) are related with each other by $T_K^{\rm SE} = T_K^2/\sqrt{\Gamma U}$, which implies that the scale $\Delta/T_K^{\rm SE}$ differs from the scale given in our work [2]. The unusual definition of Kondo temperature in Eq. (1) explains the (otherwise) unusual behaviors of $I(\phi)$ with respect to U/Δ in Fig. 2 of SE.

We now move on to point (ii). In Ref. [1] all calculations have been done at a finite temperature $T=0.1\Delta$ and SE note that "this appears to be quite close to the ground-state limit". This is particularly important in the

determination of the current-phase relation. To estimate the Josephson energy we note that it is obtained from

$$E_J(\phi) = \int^{\phi} d\phi' \ I_S(\phi') \sim \Delta \frac{I_c}{I_c^{\text{short}}}, \tag{4}$$

where I_c is the effective critical current of the system and $I_c^{\rm short} \equiv e\Delta/\hbar$ the critical current of the open contact. According the the numerical results in Ref. [1], $I_c/I_c^{\rm short} \leq 0.1$ for $\Delta/T_K^{\rm SE} \gtrsim 5$ ($\Delta/T_K \gtrsim 1$ in Ref. [2]). We think that in most plots in Ref. [1] the current-phase relation contains significant amount of thermal activation. To confirm this we have performed NRG calculations at finite temperatures and the results in Fig. 1 demonstrate the strong finite-temperature effects. The sharp transition at zero temperature is washed out and the critical current is reduced by a factor of 5 for $T/\Delta=0.1$. The discrepancy between the NRG and QMC data in the new Fig. 2 of the Reply[7] may simply reflect the different estimates of critical value Δ_c/T_K (i.e., the NRG and QMC data are in different phases), and may not be an evidence that the NRG is less accurate.

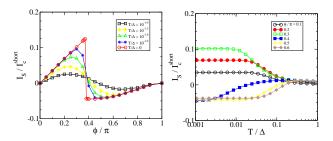


FIG. 1: (a) Josephson current $I_S(\phi)$ at different temperatures. (b) Josephson current as a function of temperature for different values of ϕ . $\Delta/T_K=1.6$.

Authors:

Mahn-Soo Choi and Minchul Lee, Department of Physics, Korea University, Seoul 136-701, Korea Kicheon Kang, Department of Physics, Chonnam National University, Gwang-ju 500-757, Korea W. Belzig, University of Basel, Klingelbergstrasse 82, 4056 Basel, Switzerland

^[1] F. Siano and R. Egger, Phys. Rev. Lett. 93, 047002 (2004).

^[2] M.-S. Choi, M. Lee, K. Kang, and W. Belzig, Phys. Rev. B 70, R020502 (2004).

^[3] T. Yoshioka and Y. Ohashi, J. Phys. Soc. Jpn. 69, 1812 (2000).

^[4] F. D. M. Haldane, Phys. Rev. Lett. 40, 416 (1978); 911 (1978).

^[5] A. C. Hewson, The Kondo Problem to Heavy Fermions (Cambridge 1993).

- [6] W. G. van der Wiel et al., Science $\bf 289,\ 2105$ (2000). [7] F. Siano and R. Egger, cond-mat/0410462 (unpublished).